

LIQUID CRYSTALS AND LIGHT EMITTING MATERIALS FOR PHOTONIC APPLICATIONS

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Lecture series at WAT in Warsaw

OVERVIEW

Liquid crystal properties (10h)

Properties of nematic liquid crystals

Nematic order parameter

Polarization and dielectric constant

Elastic energy

Surface alignment

Electrical energy

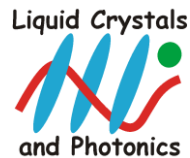
Jones matrix method

Variable phase retarder

VAN mode

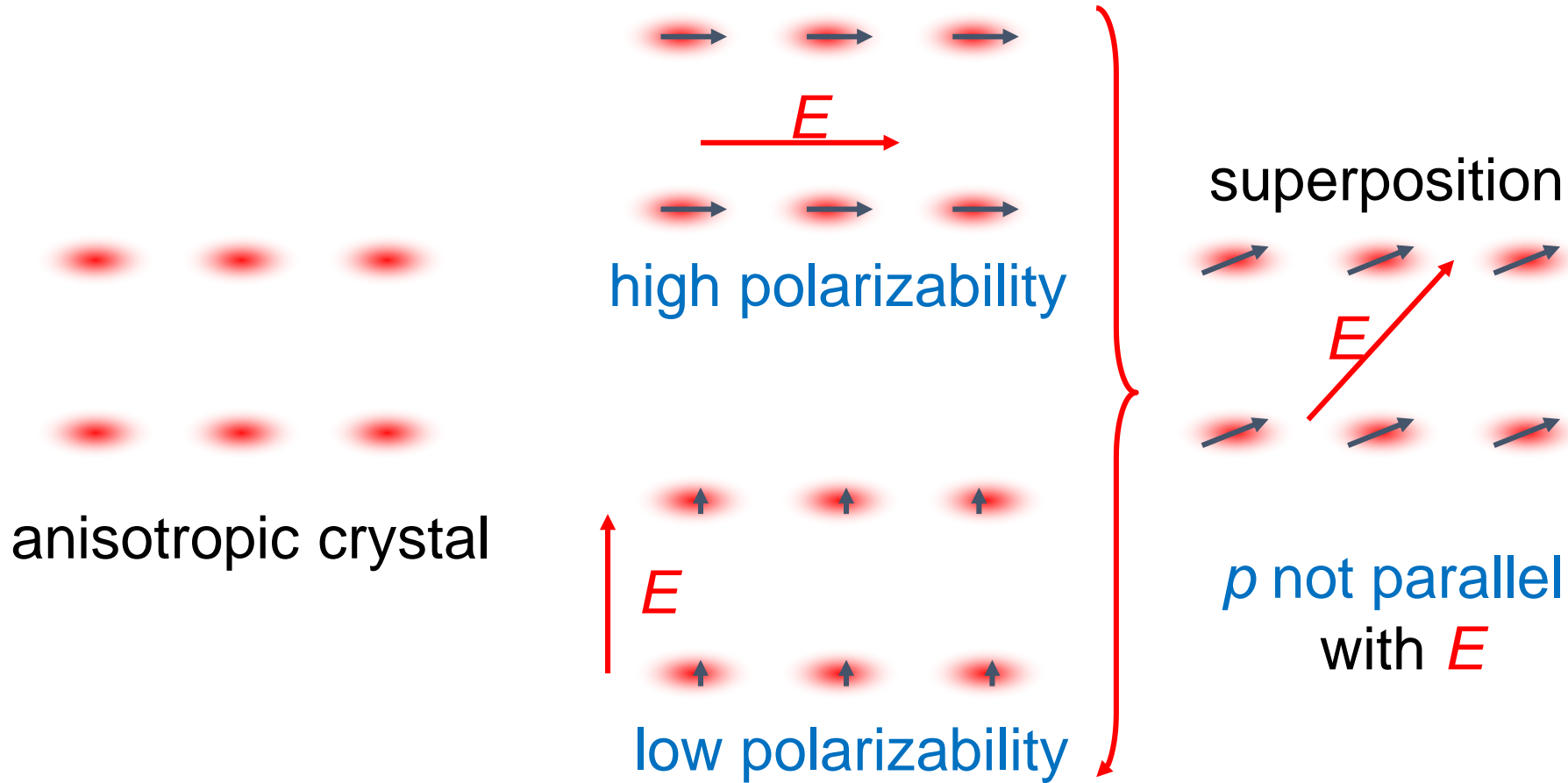
IPS mode

TN mode



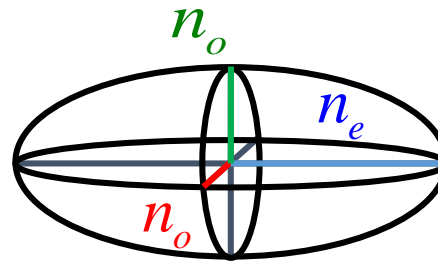
ANISOTROPIC POLARIZABILITY

Anisotropic polarizability in a crystalline structure

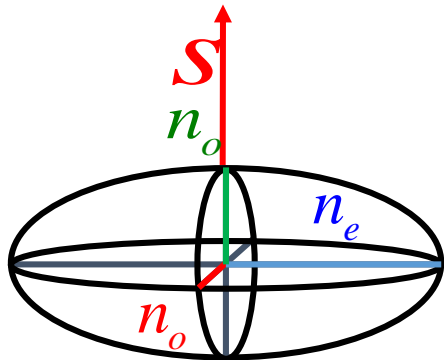


LIGHT PROPAGATION

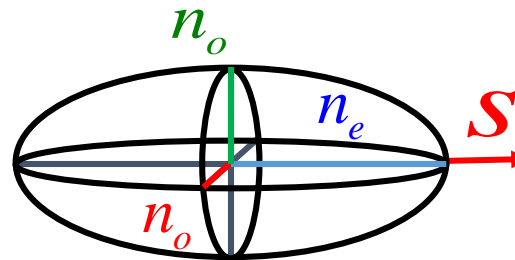
$$\frac{X^2}{n_o^2} + \frac{Y^2}{n_o^2} + \frac{Z^2}{n_e^2} = 1$$



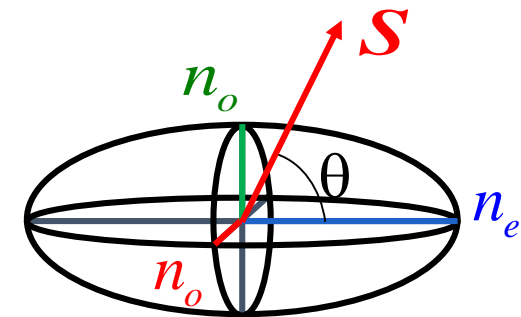
$$\bar{\bar{\eta}} = \begin{pmatrix} \frac{1}{n_o^2} & 0 & 0 \\ 0 & \frac{1}{n_o^2} & 0 \\ 0 & 0 & \frac{1}{n_e^2} \end{pmatrix}$$



two plane waves
with n_o and n_e



two plane waves
with n_o and n_o



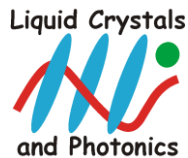
two plane waves
with n_o and

$$\frac{1}{\sqrt{\frac{\cos^2 \theta}{n_o^2} + \frac{\sin^2 \theta}{n_e^2}}}$$

HOMOGENEOUS NEMATIC LIQUID CRYSTAL

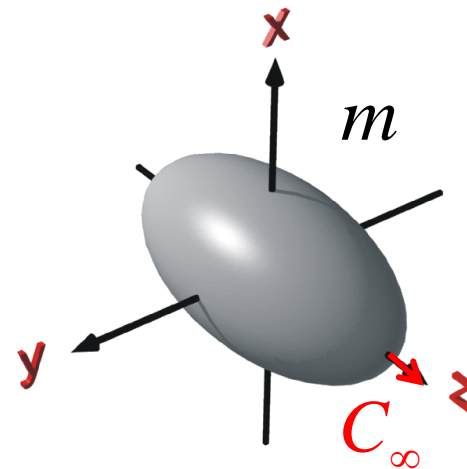
Symmetry elements for nematic LC

- full rotation symmetry axis C_∞ (z-axis)
- mirror plane m (xy plane)
- C_2 axes in the plane xy
- mirror planes containing the z-axis



tensors of rank 2

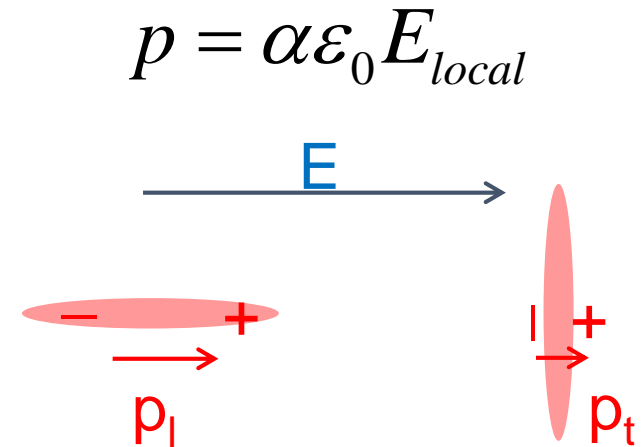
have **uniaxial** symmetry



LIQUID CRYSTAL POLARIZABILITY

Anisotropic polarization:

polarization due to
molecular polarizability α
induced dipoles



$$p = \alpha \varepsilon_0 E_{local}$$

$$p_l = \alpha_l \varepsilon_0 E_{local,l}$$

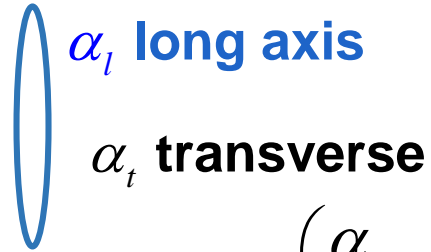
$$p_t = \alpha_t \varepsilon_0 E_{local,t}$$

$$\Delta \alpha = \alpha_l - \alpha_t > 0$$

$$\alpha_{avg} = \frac{1}{3} \alpha_l + \frac{2}{3} \alpha_t$$

MICROSCOPIC THEORY FOR UNIAXIAL DIELECTRICS

Polarizability tensor



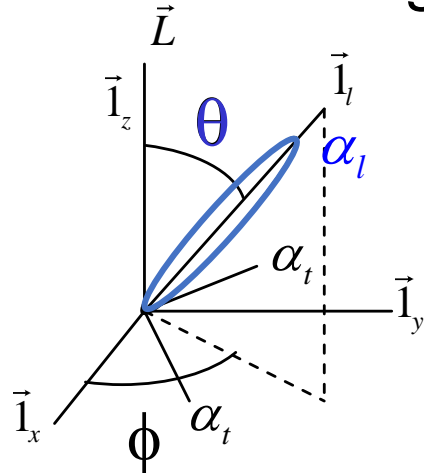
for a molecule, related to its axes

α_t is average of two transverse α 's
average of three α 's: $\alpha_{avg} = \frac{\alpha_l + 2\alpha_t}{3}$

$$\bar{\bar{\alpha}}_{mol} = \begin{pmatrix} \alpha_t & 0 & 0 \\ 0 & \alpha_t & 0 \\ 0 & 0 & \alpha_l \end{pmatrix}$$

for a molecule, related to the axes of nematic phase

molecule along \vec{l}_l has inclination θ and azimuth ϕ



$$\vec{l}_l = \begin{pmatrix} \cos \phi \sin \theta \\ \sin \phi \sin \theta \\ \cos \theta \end{pmatrix}$$

$\vec{l}_l \neq \vec{L}$ due to thermal motion

MICROSCOPIC THEORY FOR UNIAXIAL DIELECTRICS

Coordinate transformations

OLD to INTERMEDIATE

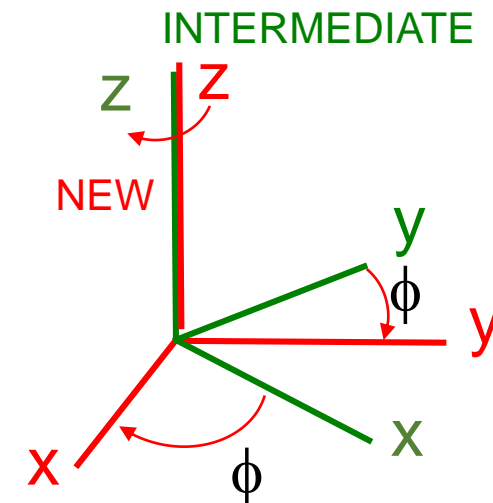
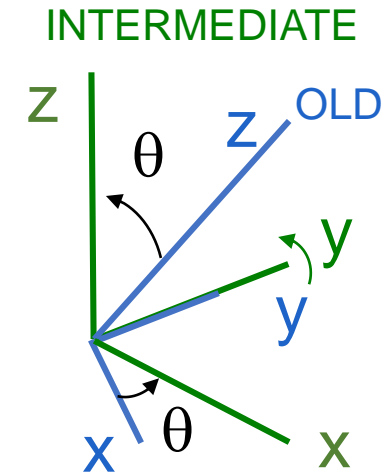
$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} \cos \theta & 0 & \sin \theta \\ 0 & 1 & 0 \\ -\sin \theta & 0 & \cos \theta \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

INTERMEDIATE to NEW

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} \cos \phi & -\sin \phi & 0 \\ \sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

columns are xyz coordinates of old unit vectors

$$\begin{pmatrix} \cos \phi \cos \theta & -\sin \phi & \cos \phi \sin \theta \\ \sin \phi \cos \theta & \cos \phi & \sin \phi \sin \theta \\ -\sin \theta & 0 & \cos \theta \end{pmatrix}$$



TENSOR TRANSFORMATION FORMULA

Transformation of a tensor?

transformation of p (old coordinates)

to a **new system** of axes p' (**new coordinates**)

$$p_k' = R_{ki} p_i$$

$$p_i = \varepsilon_0 \alpha_{ij} E_j$$

$$p_k' = R_{ki} p_i$$

$$E_l' = R_{lj} E_j \Rightarrow E_j = R_{lj} E_l'$$

$$p_k' = \varepsilon_0 \underline{\alpha_{kl}'} E_l'$$

$$p_k' = \varepsilon_0 \underline{R_{ki} \alpha_{ij} R_{lj}} E_l'$$

$$\alpha_{kl}' = R_{ki} R_{lj} \alpha_{ij}$$

TENSOR TRANSFORMATION FORMULA

Transformation of a tensor?

transformation of α (old coordinates)

to a **new system** of axes α' (**new coordinates**)

$$\alpha_{kl}' = R_{ki} R_{lj} \alpha_{ij}$$

$$\bar{\bar{\alpha}}_{mol} = \begin{pmatrix} \alpha_t & 0 & 0 \\ 0 & \alpha_t & 0 \\ 0 & 0 & \alpha_l \end{pmatrix}$$

$$\bar{\bar{\alpha}}_{mol}' = \begin{pmatrix} \cos \phi \cos \theta & -\sin \phi & \cos \phi \sin \theta \\ \sin \phi \cos \theta & \cos \phi & \sin \phi \sin \theta \\ -\sin \theta & 0 & \cos \theta \end{pmatrix} \begin{pmatrix} \alpha_t & 0 & 0 \\ 0 & \alpha_t & 0 \\ 0 & 0 & \alpha_l \end{pmatrix} \begin{pmatrix} \cos \phi \cos \theta & -\sin \phi & \cos \phi \sin \theta \\ \sin \phi \cos \theta & \cos \phi & \sin \phi \sin \theta \\ -\sin \theta & 0 & \cos \theta \end{pmatrix}^{-1}$$

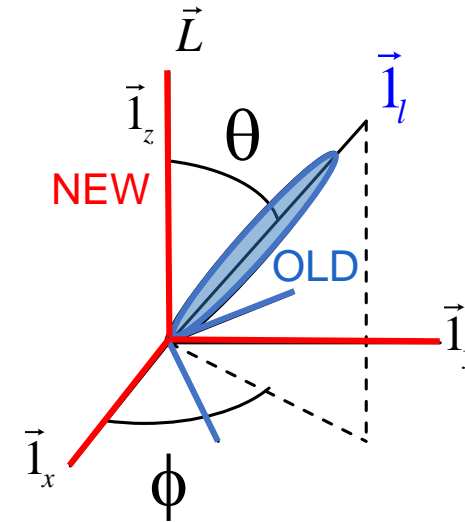
MICROSCOPIC THEORY FOR UNIAXIAL DIELECTRICS

Coordinate transformation: OLD to NEW

$$R = \begin{pmatrix} \cos \phi \cos \theta & -\sin \phi & \cos \phi \sin \theta \\ \sin \phi \cos \theta & \cos \phi & \sin \phi \sin \theta \\ -\sin \theta & 0 & \cos \theta \end{pmatrix}$$

$$\bar{\bar{\alpha}}_{mol} = \begin{pmatrix} \alpha_t & 0 & 0 \\ 0 & \alpha_t & 0 \\ 0 & 0 & \alpha_l \end{pmatrix} = \alpha_t \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \alpha_l - \alpha_t \end{pmatrix}$$

$$\alpha_{mol,mn} = \alpha_t \delta_{mn} + (\alpha_l - \alpha_t) \delta_{m3} \delta_{n3} \quad \delta_{ij}: 1 \text{ when } i=j; 0 \text{ when } i \neq j$$



transform →

$$\begin{aligned} \alpha_{ij} &= R_{im} R_{jn} \alpha_{mol,mn} = R_{im} R_{jn} \alpha_t \delta_{mn} + (\alpha_l - \alpha_t) R_{im} R_{jn} \delta_{m3} \delta_{n3} \\ &= R_{im} R_{jm} \alpha_t + (\alpha_l - \alpha_t) R_{i3} R_{j3} \\ &= \alpha_t \delta_{ij} + (\alpha_l - \alpha_t) l_i l_j \end{aligned}$$

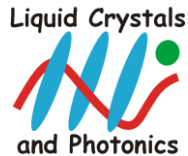
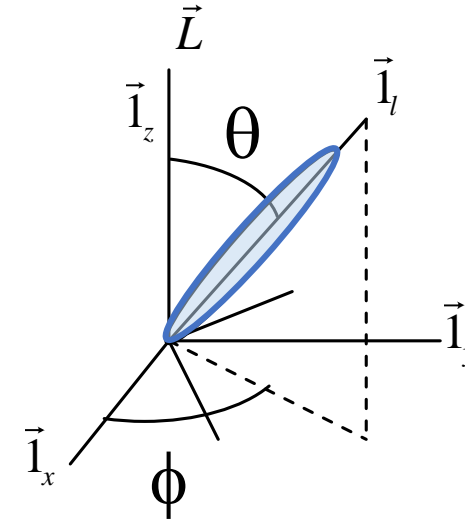
MICROSCOPIC THEORY FOR UNIAXIAL DIELECTRICS

Value of α in xyz ?

$$\alpha_{ij} = \alpha_t \delta_{ij} + (\alpha_l - \alpha_t) l_i l_j$$

$$\bar{\bar{\alpha}} = \alpha_t \bar{\bar{\delta}} + (\alpha_l - \alpha_t) \bar{\bar{l}}$$

diade product: $\bar{\bar{l}} = \begin{pmatrix} l_x \\ l_y \\ l_z \end{pmatrix} \begin{pmatrix} l_x & l_y & l_z \end{pmatrix}$



$$\bar{\bar{\alpha}} = \begin{pmatrix} \alpha_t + (\alpha_l - \alpha_t) \cos^2 \phi \sin^2 \theta & (\alpha_l - \alpha_t) \sin \phi \cos \phi \sin^2 \theta & (\alpha_l - \alpha_t) \cos \phi \sin \theta \cos \theta \\ (\alpha_l - \alpha_t) \sin \phi \cos \phi \sin^2 \theta & \alpha_t + (\alpha_l - \alpha_t) \sin^2 \phi \sin^2 \theta & (\alpha_l - \alpha_t) \sin \phi \sin \theta \cos \theta \\ (\alpha_l - \alpha_t) \cos \phi \sin \theta \cos \theta & (\alpha_l - \alpha_t) \sin \phi \sin \theta \cos \theta & \alpha_t + (\alpha_l - \alpha_t) \cos^2 \theta \end{pmatrix}$$



$\langle \bar{\bar{\alpha}} \rangle$? average over orientations

ϕ random between 0 and 2π

$$\begin{cases} \langle \cos \phi \rangle = \langle \sin \phi \rangle = \langle \cos \phi \sin \phi \rangle = 0 \\ \langle \cos^2 \phi \rangle = \frac{1}{2} \end{cases}$$



MICROSCOPIC THEORY FOR UNIAXIAL DIELECTRICS

average polarizability tensor

$$\langle \bar{\alpha} \rangle = \begin{pmatrix} \alpha_t + \frac{1}{2}(\alpha_l - \alpha_t) \langle \sin^2 \theta \rangle & 0 & 0 \\ 0 & \alpha_t + \frac{1}{2}(\alpha_l - \alpha_t) \langle \sin^2 \theta \rangle & 0 \\ 0 & 0 & \alpha_t + (\alpha_l - \alpha_t) \langle \cos^2 \theta \rangle \end{pmatrix}$$

definition for nematic phase: \perp and \parallel with director \vec{L}

$$\langle \bar{\alpha} \rangle = \begin{pmatrix} \alpha_{\perp} & 0 & 0 \\ 0 & \alpha_{\perp} & 0 \\ 0 & 0 & \alpha_{\parallel} \end{pmatrix} \Rightarrow \begin{cases} \alpha_{\perp} = \alpha_t + \frac{1}{2}(\alpha_l - \alpha_t) \langle \sin^2 \theta \rangle \\ \alpha_{\parallel} = \alpha_t + (\alpha_l - \alpha_t) \langle \cos^2 \theta \rangle \end{cases}$$

for a molecule:

$$\alpha_{avg} = \frac{\alpha_l + 2\alpha_t}{3}$$

check that:

$$\frac{\alpha_{\parallel} + 2\alpha_{\perp}}{3} = \alpha_{avg}$$

MICROSCOPIC THEORY FOR UNIAXIAL DIELECTRICS

Relation with the order parameter

$$S = \frac{1}{2} \langle 3 \cos^2 \theta - 1 \rangle$$

$$\langle \cos^2 \theta \rangle = \frac{1 + 2S}{3}$$

$$\langle \sin^2 \theta \rangle = \frac{2 - 2S}{3}$$



$$\alpha_{//} = \alpha_t + (\alpha_l - \alpha_t) \langle \cos^2 \theta \rangle = \dots = \alpha_{avg} + (\alpha_l - \alpha_t) \frac{2}{3} S$$

$$\alpha_{\perp} = \alpha_t + \frac{1}{2} (\alpha_l - \alpha_t) \langle \sin^2 \theta \rangle = \dots = \alpha_{avg} - (\alpha_l - \alpha_t) \frac{1}{3} S$$

check that: $\alpha_{//} - \alpha_{\perp} = (\alpha_l - \alpha_t) S$

MICROSCOPIC THEORY FOR UNIAXIAL DIELECTRICS

$$\langle \bar{\alpha} \rangle = \begin{pmatrix} \alpha_{\perp} & 0 & 0 \\ 0 & \alpha_{\perp} & 0 \\ 0 & 0 & \alpha_{\parallel} \end{pmatrix} = \begin{pmatrix} \alpha_{avg} - (\alpha_l - \alpha_t) \frac{1}{3} S & 0 & 0 \\ 0 & \alpha_{avg} - (\alpha_l - \alpha_t) \frac{1}{3} S & 0 \\ 0 & 0 & \alpha_{avg} + (\alpha_l - \alpha_t) \frac{2}{3} S \end{pmatrix}$$

Susceptibility tensor?

$$\bar{\chi} \approx N \langle \bar{\alpha} \rangle$$

$$\chi_{\parallel} \approx N \alpha_{\parallel}$$

$$\chi_{\perp} \approx N \alpha_{\perp}$$

$$\bar{\varepsilon} = \varepsilon_0 (1 + \bar{\chi})$$

DIELECTRIC TENSOR

Dielectric tensor? from D -field

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P} = \epsilon_0 (1 + \overline{\chi}) \vec{E} = \overline{\epsilon} \vec{E} \quad \overline{\epsilon} = \begin{pmatrix} \epsilon_{\perp} & 0 & 0 \\ 0 & \epsilon_{\perp} & 0 \\ 0 & 0 & \epsilon_{\parallel} \end{pmatrix}$$

$$\epsilon_{\parallel} = \epsilon_0 (1 + \chi_{\parallel})$$

$$\epsilon_{\perp} = \epsilon_0 (1 + \chi_{\perp})$$

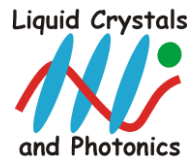
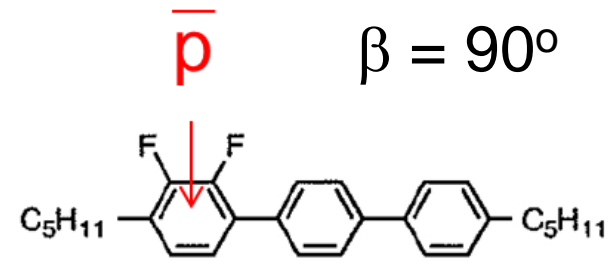
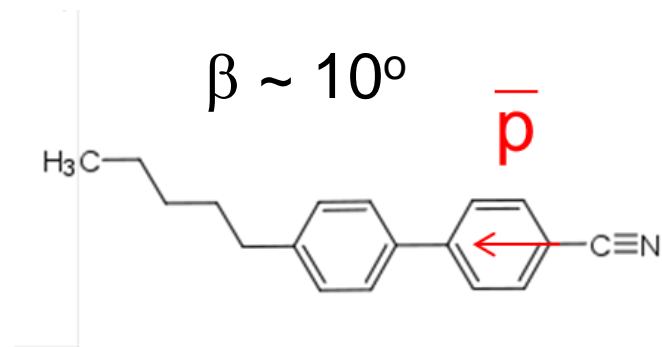
$$\epsilon_{avg} = \frac{\epsilon_{\parallel} + 2\epsilon_{\perp}}{3} \quad \downarrow \quad \Rightarrow \quad \frac{\epsilon_{\parallel} - \epsilon_{\perp}}{\epsilon_{avg} - \epsilon_0} = \frac{\chi_{\parallel} - \chi_{\perp}}{\chi_{avg}}$$

Conclusion

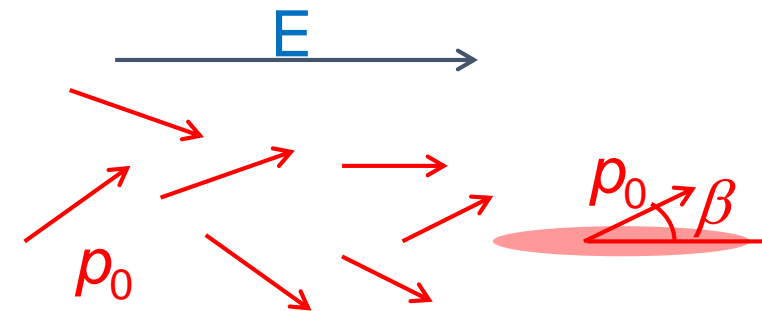
$$\frac{\epsilon_{\parallel} - \epsilon_{\perp}}{\epsilon_{avg} - \epsilon_0} = \frac{\chi_{\parallel} - \chi_{\perp}}{\chi_{avg}} = \frac{\alpha_{\parallel} - \alpha_{\perp}}{\alpha_{avg}} = \frac{\alpha_l - \alpha_t}{\alpha_{avg}} S$$

LIQUID CRYSTAL POLARIZABILITY

molecules with permanent dipole moment

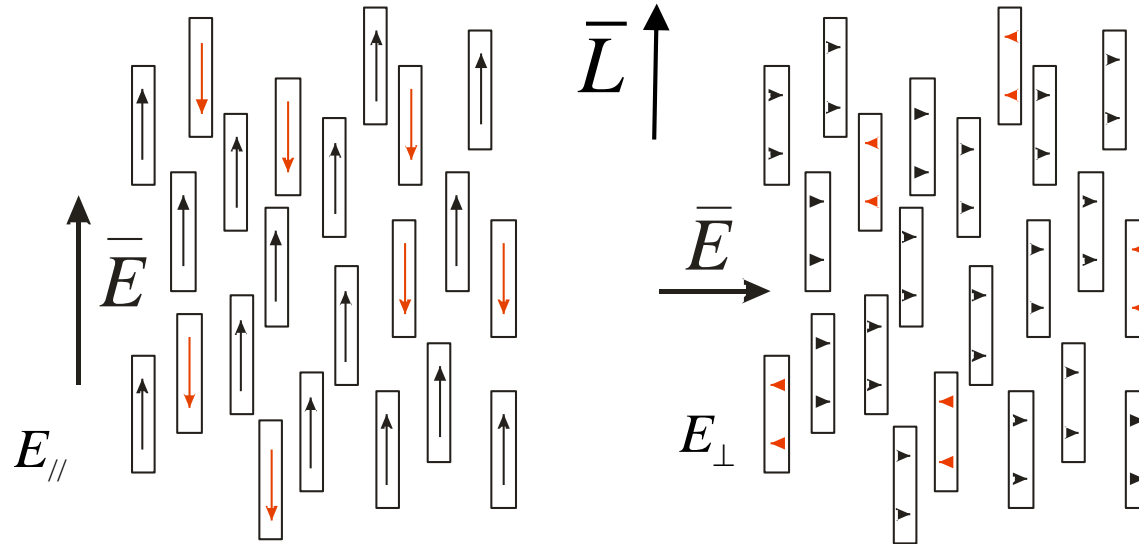


orientation polarization
permanent dipoles
(low frequencies)



ELECTRICAL PROPERTIES

More dipoles parallel to the electric field (Boltzmann, Maier)



shown for $S=1$
(perfect order)

$\beta = 0^\circ$

$\beta = 90^\circ$

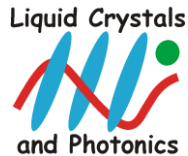
h : cavity field factor
 F : reaction field factor

$$P_{ori, //} = \frac{F^2 p_0^2}{3kT} [1 + 2S] h E_{//} \quad [1 - (1 - 3\cos^2 \beta)S]$$

$$P_{ori, //} = \frac{F^2 p_0^2}{3kT} [1 - S] h E_{//} \quad // \text{ long axis}$$

$$P_{ori, \perp} = \frac{F^2 p_0^2}{3kT} [1 - S] h E_{\perp} \quad [1 + \frac{1}{2}(1 - 3\cos^2 \beta)S]$$

$$P_{ori, \perp} = \frac{F^2 p_0^2}{3kT} \left[1 + \frac{1}{2} S \right] h E_{\perp} \quad \perp \text{ long axis}$$



ELECTRICAL PROPERTIES

Anisotropic (uniaxial) dielectric :

$$\begin{pmatrix} D_x \\ D_y \\ D_z \end{pmatrix} = \epsilon_0 \begin{pmatrix} \epsilon_{\perp} & 0 & 0 \\ 0 & \epsilon_{\perp} & 0 \\ 0 & 0 & \epsilon_{\parallel} \end{pmatrix} \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix}$$

h : cavity field factor

F : reaction field factor

polarization due to
molecular polarizability α
induced dipoles

$$\begin{cases} p_{ele,\parallel} = \epsilon_0 F \left(\alpha_{avg} + \frac{2}{3} \Delta\alpha S \right) h E_{\parallel} \\ p_{ele,\perp} = \epsilon_0 F \left(\alpha_{avg} - \frac{1}{3} \Delta\alpha S \right) h E_{\perp} \end{cases}$$

orientation polarization
permanent dipoles
(low frequencies)

$$\begin{cases} p_{ori,\parallel} = \frac{F^2 p_0^2}{3kT} \left[1 - (1 - 3 \cos^2 \beta) S \right] h E_{\parallel} \\ p_{ori,\perp} = \frac{F^2 p_0^2}{3kT} \left[1 + \frac{1}{2} (1 - 3 \cos^2 \beta) S \right] h E_{\perp} \end{cases}$$

DIELECTRIC PROPERTIES

Anisotropic (uniaxial) dielectric :

$$\begin{pmatrix} D_x \\ D_y \\ D_z \end{pmatrix} = \epsilon_0 \begin{pmatrix} \epsilon_{\perp} & 0 & 0 \\ 0 & \epsilon_{\perp} & 0 \\ 0 & 0 & \epsilon_{\parallel} \end{pmatrix} \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix}$$

$$\epsilon_{\parallel} = \epsilon_0 + \epsilon_0 N F h \left(\alpha_{avg} + \frac{2}{3} \Delta \alpha S + \frac{F p_0^2}{3 k T} \left[1 - (1 - 3 \cos^2 \beta) S \right] \right)$$

$$\epsilon_{\perp} = \epsilon_0 + \epsilon_0 N F h \left(\alpha_{avg} - \frac{1}{3} \Delta \alpha S + \frac{F p_0^2}{3 k T} \left[1 + \frac{1}{2} (1 - 3 \cos^2 \beta) S \right] \right)$$

$$\Delta \epsilon = \epsilon_{\parallel} - \epsilon_{\perp}$$

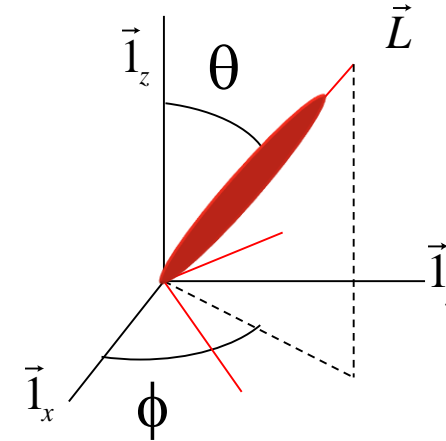
only at low frequencies <100kHz
orientation polarization is slow

$\Delta \epsilon < 0$ at low frequency if $\beta = 90^\circ$

DIELECTRIC TENSOR

Dielectric tensor in laboratory axes?

$$\bar{\epsilon}_{\parallel} = \begin{pmatrix} \epsilon_{\perp} & 0 & 0 \\ 0 & \epsilon_{\perp} & 0 \\ 0 & 0 & \epsilon_{\parallel} \end{pmatrix} \quad \vec{L} = \begin{pmatrix} \cos \phi \sin \theta \\ \sin \phi \sin \theta \\ \cos \theta \end{pmatrix}$$



$$R = \begin{pmatrix} \cos \phi \cos \theta & -\sin \phi & \cos \phi \sin \theta \\ \sin \phi \cos \theta & \cos \phi & \sin \phi \sin \theta \\ -\sin \theta & 0 & \cos \theta \end{pmatrix}$$

$$\begin{aligned} \epsilon_{\alpha\beta} &= R_{\alpha k} R_{\beta l} \epsilon_{kl} \\ &= \epsilon_{\perp} \delta_{\alpha\beta} + (\epsilon_{\parallel} - \epsilon_{\perp}) R_{\alpha 3} R_{\beta 3} \\ &= \epsilon_{\perp} \delta_{\alpha\beta} + (\epsilon_{\parallel} - \epsilon_{\perp}) L_{\alpha} L_{\beta} \end{aligned} \quad \epsilon_{kl} = \epsilon_{\perp} \delta_{kl} + (\epsilon_{\parallel} - \epsilon_{\perp}) \delta_{k3} \delta_{l3}$$

DIELECTRIC TENSOR

Dielectric tensor

$$\overline{\overline{\epsilon}} = \begin{pmatrix} \epsilon_{\perp} + (\epsilon_{//} - \epsilon_{\perp}) \cos^2 \phi \sin^2 \theta & (\epsilon_{//} - \epsilon_{\perp}) \sin \phi \cos \phi \sin^2 \theta & (\epsilon_{//} - \epsilon_{\perp}) \cos \phi \sin \theta \cos \theta \\ (\epsilon_{//} - \epsilon_{\perp}) \sin \phi \cos \phi \sin^2 \theta & \epsilon_{\perp} + (\epsilon_{//} - \epsilon_{\perp}) \sin^2 \phi \sin^2 \theta & (\epsilon_{//} - \epsilon_{\perp}) \sin \phi \sin \theta \cos \theta \\ (\epsilon_{//} - \epsilon_{\perp}) \cos \phi \sin \theta \cos \theta & (\epsilon_{//} - \epsilon_{\perp}) \sin \phi \sin \theta \cos \theta & \epsilon_{\perp} + (\epsilon_{//} - \epsilon_{\perp}) \cos^2 \theta \end{pmatrix}$$

equivalent with:

$$\epsilon_{\alpha\beta} = \epsilon_{\perp} \delta_{\alpha\beta} + (\epsilon_{//} - \epsilon_{\perp}) L_{\alpha} L_{\beta}$$

$$\overline{\overline{\epsilon}} = \epsilon_{\perp} \overline{\overline{\delta}} + (\epsilon_{//} - \epsilon_{\perp}) \vec{L} \vec{L}$$

$$= \epsilon_{\perp} \overline{\overline{\delta}} + \Delta\epsilon \vec{L} \vec{L}$$

$$\vec{L} = \begin{pmatrix} \cos \phi \sin \theta \\ \sin \phi \sin \theta \\ \cos \theta \end{pmatrix}$$

diade product

$$\vec{L} \vec{L} = \begin{pmatrix} L_x \\ L_y \\ L_z \end{pmatrix} \begin{pmatrix} L_x & L_y & L_z \end{pmatrix}$$

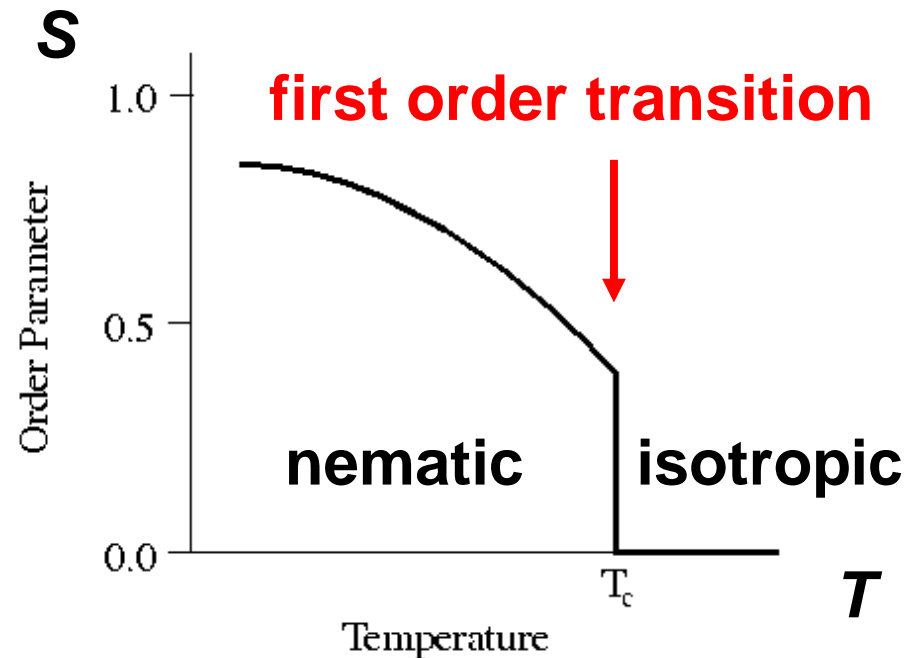
INFLUENCE OF TEMPERATURE

Transition from nematic to isotropic

is related with a jump in the order parameter

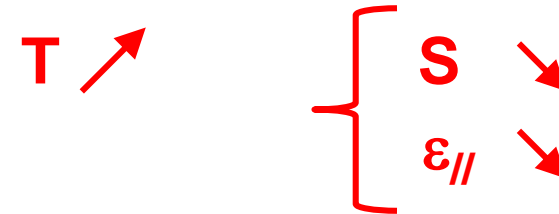
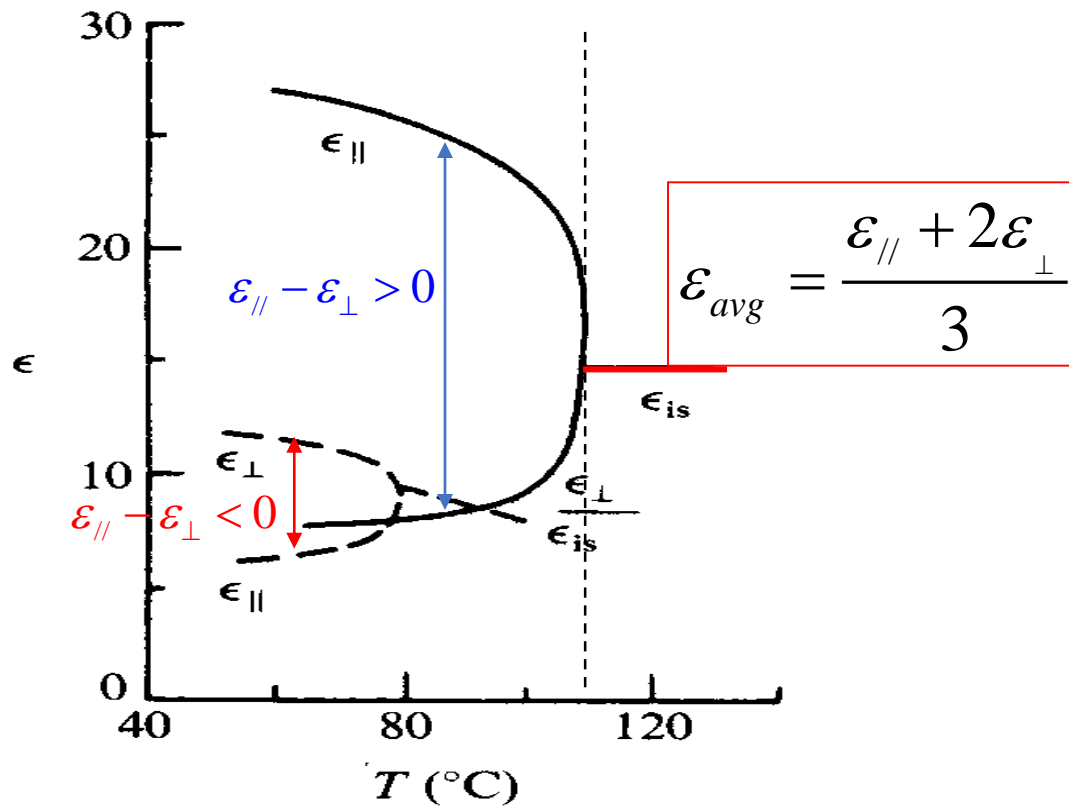
this is called a **first order transition**

(with change in enthalpy, heat is needed to cross T_c)



DIELECTRIC TENSOR

nematic isotropic



low frequencies

$$\left[\begin{array}{l} 5\epsilon_0 < \epsilon_{||}, \epsilon_{\perp} < 80\epsilon_0 \\ -5\epsilon_0 < \epsilon_{||} - \epsilon_{\perp} < 20\epsilon_0 \end{array} \right.$$

optical frequencies

$$\left[\begin{array}{l} 2\epsilon_0 < \epsilon_{||}, \epsilon_{\perp} < 4\epsilon_0 \\ 0 < \epsilon_{||} - \epsilon_{\perp} < 1.5\epsilon_0 \end{array} \right.$$

REFRACTIVE INDICES

Ordinary and extra-ordinary indices:

$$n_e = \sqrt{\frac{\epsilon_{//}}{\epsilon_0}}; \quad n_o = \sqrt{\frac{\epsilon_{\perp}}{\epsilon_0}}$$

optical frequencies: $\left\{ \begin{array}{l} 1.4 < n_e, n_o < 2 \\ 0 < n_e - n_o < 0.4 \end{array} \right.$

For light propagation with k -vector at an angle θ from L

$$n^{(2)} = \left(\frac{\cos^2 \theta}{n_o^2} + \frac{\sin^2 \theta}{n_e^2} \right)^{-\frac{1}{2}}$$

TENSOR OF MAGNETIC PERMEABILITY

Magnetic permeability tensor $\vec{M} = \mu_0 \overset{=}{\chi}_m \vec{H}$

$$\overset{=}{\chi}_m = \begin{pmatrix} \chi_{m\perp} & 0 & 0 \\ 0 & \chi_{m\perp} & 0 \\ 0 & 0 & \chi_{m\parallel} \end{pmatrix}$$

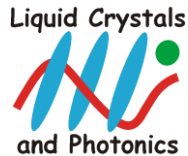
diamagnetism: $\chi < 0$ $\chi_{m\perp}, \chi_{m\parallel} < 0$

easier magnetization for the phenyl ring, perpendicular to director

$$|\chi_{m\perp}| > |\chi_{m\parallel}| \quad \Rightarrow \quad \Delta\chi_m = \chi_{m\parallel} - \chi_{m\perp} > 0$$

relation with order parameter:
($H_{\text{local}} \sim H$)

$$S = \frac{\chi_{\parallel} - \chi_{\perp}}{(\chi_{\parallel} - \chi_{\perp})_{\text{max}}}$$



INFLUENCE OF ELECTRIC FIELD

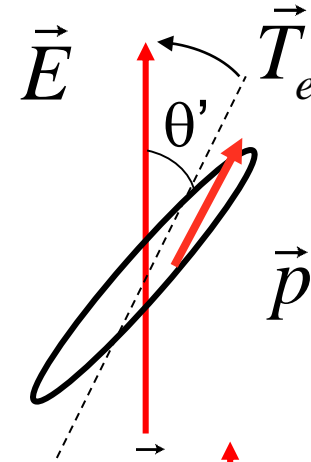
Torque by electric field on molecules:

$$\vec{T}_e = \vec{p} \times \vec{E}_{local}$$

permanent dipole moment

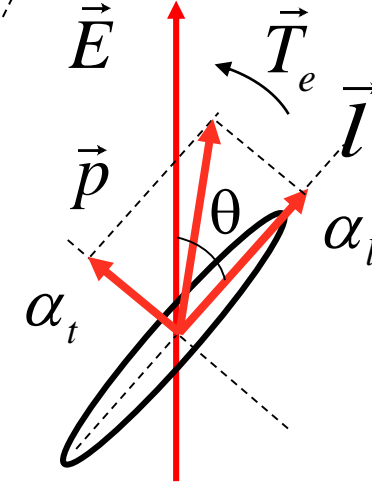
p aligns with E , due to torque T_e

$$|\vec{T}_e| = pE_{local} \sin \theta'$$



induced dipole moment $\vec{p} = \epsilon_0 \bar{\alpha} \vec{E}_{local}$

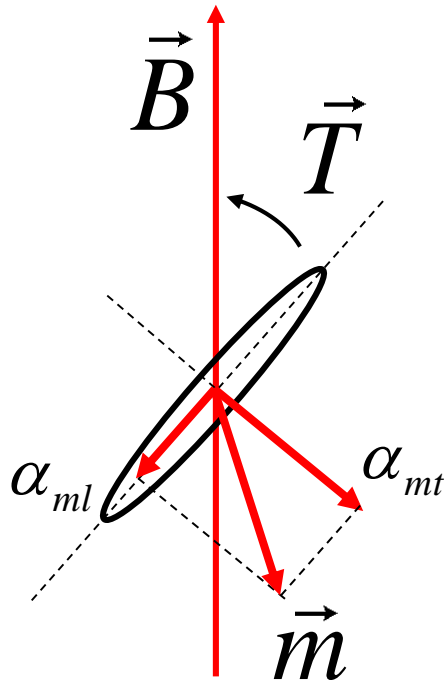
$$\begin{aligned} |\vec{T}_e| &= \left| \epsilon_0 \bar{\alpha} \vec{E}_{local} \times \vec{E}_{local} \right| \\ &= \left| \left(\epsilon_0 \alpha_t \bar{\delta} \vec{E}_{local} + \epsilon_0 (\alpha_l - \alpha_t) (\vec{l} \vec{l}) \vec{E}_{local} \right) \times \vec{E}_{local} \right| \\ &= \left| \left(\epsilon_0 (\alpha_l - \alpha_t) \cos \theta E_{local} \vec{l} \right) \times \vec{E}_{local} \right| \\ &= \epsilon_0 (\alpha_l - \alpha_t) \cos \theta \sin \theta E_{local}^2 \end{aligned}$$



for $\alpha_l > \alpha_t$

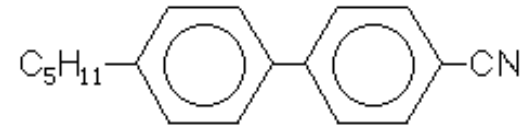
INFLUENCE OF MAGNETIC FIELD

Torque by magnetic field on molecules



diamagnetism
(stronger for phenyl rings)

$$\alpha_{mt} < \alpha_{ml} < 0$$



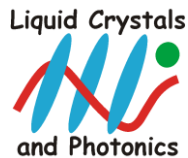
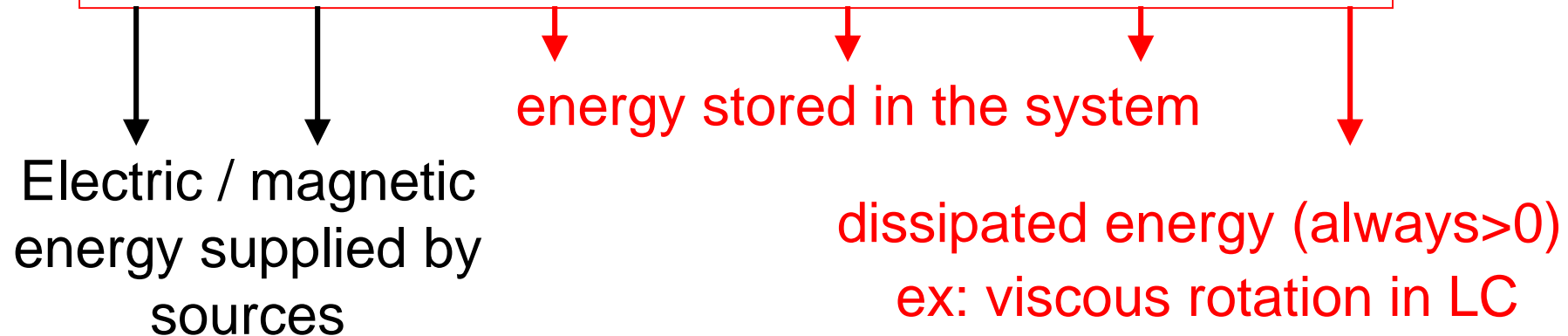
Torque $\vec{T}_m = \vec{m} \times \vec{B}$

aligns long axis with B

PRINCIPLE OF VIRTUAL WORK

Conservation of energy, in the presence of a source

$$VdQ + Id\Phi = dW_{electric} + dW_{magnetic} + dW_{elastic} + dA$$



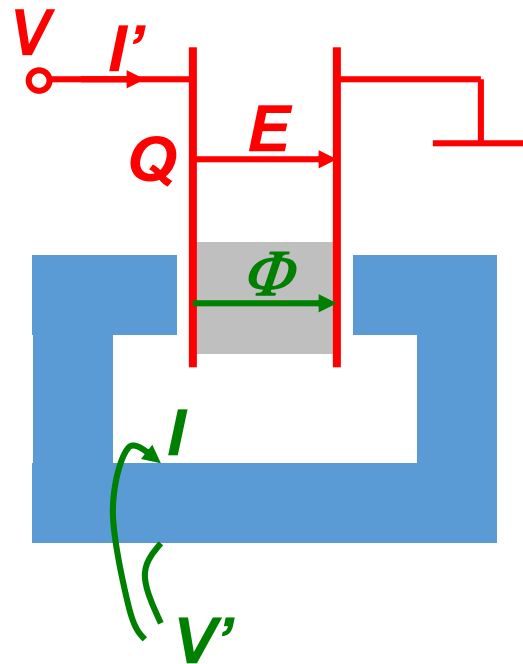
“Stable director distribution is obtained if every virtual variation leads to $dA < 0$ (not allowed)”

Without sources, the **stored energy** should be minimal

PRINCIPLE OF VIRTUAL WORK

Using sources that can supply energy to the LC

electric: source with cst V or magnetic: source with cst I



Electric

energy

$$VdQ$$

current

$$I' = \frac{dQ}{dt}$$

Magnetic

energy

$$Id\Phi$$

voltage

$$V' = \frac{d\Phi}{dt}$$

PRINCIPLE OF VIRTUAL WORK

Internal energy in the system

$$\left\{ \begin{array}{l} W_{electric} = \int_{vol} \frac{1}{2} \vec{D} \cdot \vec{E} dV = \frac{1}{2} V \cdot Q \\ W_{magnetic} = \int_{vol} \frac{1}{2} \vec{B} \cdot \vec{H} dV = \frac{1}{2} I \cdot \Phi \end{array} \right.$$

cst V over a capacitor
with charge Q

cst I over a gap
with flux Φ

virtual
work

$$VdQ + Id\Phi = d\left(\frac{1}{2} V \cdot Q\right) + d\left(\frac{1}{2} I \cdot \Phi\right) + dW_{elastic} + dA$$

$$0 = -\frac{1}{2} d(VQ) - \frac{1}{2} d(I\Phi) + dW_{elastic} + dA$$

$$0 = d \int_{vol} \left(-\frac{1}{2} \vec{D} \cdot \vec{E}\right) dV + d \int_{vol} \left(-\frac{1}{2} \vec{B} \cdot \vec{H}\right) dV + dW_{elastic} + dA$$

$$0 = d \int_{vol} \left(f_{electric} + \underbrace{f_{magnetic}} + f_{elastic} \right) dV + \underbrace{dA}_{\text{dissipation} > 0}$$

energy density

integral should be minimal for stability